

**ASYMPTOTIC EXPANSIONS FOR THE WALLIS SEQUENCE AND SOME NEW MATHEMATICAL CONSTANTS ASSOCIATED WITH THE GLAISHER-KINKELIN AND CHOI-SRIVASTAVA CONSTANTS**

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The celebrated Wallis sequence  $W_n$ , which is defined by  $W_n := \prod_{k=1}^n \frac{4k^2}{4k^2-1}$ , is known to have the limit  $\frac{\pi}{2}$  as  $n \rightarrow \infty$ . Without using the Bernoulli numbers  $B_n$ , the authors present several asymptotic expansions and a recurrence relation for determining the coefficients of each asymptotic expansion related to the Wallis sequence  $W_n$  and the newly-introduced constants  $D$  and  $E$ , which are analogous to the Glaisher-Kinkelin constant  $A$  and the Choi-Srivastava constants  $B$  and  $C$ .

**1. INTRODUCTION, MOTIVATION AND PRELIMINARIES**

The famous Wallis sequence  $W_n$ , defined by

$$W_n := \prod_{k=1}^n \frac{4k^2}{4k^2-1} \quad (n \in \mathbb{N} := \{1, 2, 3, \dots\}),$$

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2020 Mathematics Subject Classification. 11M06, 41A60, 40A05

Keywords and Phrases. Wallis' sequence, Asymptotic expansion, Glaisher-Kinkelin constant, Generalized Glaisher-Kinkelin constant, Barnes  $G$ -function, Choi-Srivastava constants, Stirling's formula, Riemann zeta function, Euler-Maclaurin summation formula.

has the limit value

$$W_\infty := \lim_{n \rightarrow \infty} W_n = \prod_{k=1}^{\infty} \frac{4k^2}{4k^2 - 1} = \frac{\pi}{2},$$

which was established by Wallis in 1655 (see also [12, p. 68]). Several elementary proofs of this well-known result can be found in, for example, [3, 48, 69].

An interesting geometric construction that produces the above limit value can be found in the work of Myerson [56]. Many formulas exist for the representation of  $\pi$ , and a collection of these formulas is listed in [63, 64]. For a history of  $\pi$ , see [2, 10, 12, 33].

Some inequalities and asymptotic formulas associated with the Wallis sequence  $W_n$  can be found in [13, 21, 32, 35, 40, 44, 45, 46, 47, 49, 50, 51, 52, 53, 55, 58]. For example, Deng *et al.* [32] proved, for all  $n \in \mathbb{N}$ , that

$$\frac{\pi}{2} \left( 1 - \frac{1}{4n + \alpha} \right) < W_n \leq \frac{\pi}{2} \left( 1 - \frac{1}{4n + \beta} \right)$$

with the best possible constants  $\alpha$  and  $\beta$  given by

$$\alpha = \frac{5}{2} \quad \text{and} \quad \beta = \frac{32 - 9\pi}{3\pi - 8} = 2.614909986 \dots,$$

respectively.

Chen and Paris [21] showed that the following asymptotic expansion holds true for the Wallis sequence  $W_n$ :

$$\begin{aligned} W_n &\sim \frac{\pi}{2} \exp \left( \sum_{j=1}^{\infty} \frac{\nu_j}{n^j} \right) \\ (1) \quad &= \frac{\pi}{2} \exp \left( -\frac{1}{4n} + \frac{1}{8n^2} - \frac{5}{96n^3} + \frac{1}{64n^4} - \frac{1}{320n^5} + \frac{1}{384n^6} - \frac{25}{7168n^7} + \dots \right), \end{aligned}$$

with the coefficients  $\nu_j$  given by

$$(2) \quad \nu_j = \frac{(-1)^{j+1} \left( (4 - 2^{1-j}) B_{j+1} - (j+1) \cdot 2^{-j} \right)}{j(j+1)} \quad (j \geq 1),$$

where  $B_n$  ( $n \in \mathbb{N}_0 := \mathbb{N} \cup \{0\}$ ) denote the Bernoulli numbers defined by the following generating function:

$$\frac{z}{e^z - 1} = \sum_{n=0}^{\infty} B_n \frac{z^n}{n!} \quad (|z| < 2\pi).$$

By applying Lemma 3 in [19], Chen and Paris [21] deduced the following asymptotic expansion from (1):

$$(3) \quad W_n \sim \frac{\pi}{2} \sum_{j=0}^{\infty} \frac{\mu_j}{n^j},$$

where the coefficients  $\mu_j$  are given by the recurrence relation:

$$\mu_0 = 1, \quad \text{and} \quad \mu_j = \frac{1}{j} \sum_{k=1}^j k \nu_k \mu_{j-k} \quad (j \geq 1)$$

and the coefficients  $\nu_j$  are given in (2). This produces the expansion in the inverse powers of  $n$  given by

$$W_n \sim \frac{\pi}{2} \left( 1 - \frac{1}{4n} + \frac{5}{32n^2} - \frac{11}{128n^3} + \frac{83}{2048n^4} - \frac{143}{8192n^5} + \frac{625}{65536n^6} - \frac{1843}{262144n^7} + \dots \right) \quad (n \rightarrow \infty).$$

The main object of the present paper is to first provide a recurrence relation for determining the coefficients of  $n^{-j}$  in the expansion (1) without the help of the Bernoulli numbers  $B_n$ . We also derive a recurrence relation for determining the coefficients of  $n^{-j}$  in the expansion (3) without using the coefficients  $\nu_j$ .

The double gamma function  $\Gamma_2$  and the multiple gamma functions  $\Gamma_n$  were introduced and investigated by Barnes in a series of papers [4, 5, 6, 7]. In fact, Barnes applied these functions in the theories of elliptic functions and theta functions. Nonetheless, except possibly for the citations of  $\Gamma_2$  in the exercises by Whittaker and Watson [70, p. 264] and also by Gradshteyn and Ryzhik [38, p. 661, Entry 6.441 (4); p. 937, Entry 8.333], these functions were revived only in about the middle of the 1980s in the study of the determinants of the Laplacians on the  $n$ -dimensional unit sphere  $S^n$  (see, for example, [26, 43, 57, 61, 67, 68]). The theory of the double gamma function  $\Gamma_2$  has indeed found interesting applications in many other recent investigations (see, for details, [65, 66]).

Barnes [4] defined the double gamma function (or the Barnes  $G$ -function)  $\Gamma_2 = 1/G$ , which satisfies each of the following properties:

- (i)  $G(z+1) = \Gamma(z)G(z)$  for all complex  $z$ ;
- (ii)  $G(1) = 1$ ;
- (iii) As  $n \rightarrow \infty$ ,

$$\begin{aligned} \ln G(z+n+2) = & \frac{n+1+z}{2} \ln(2\pi) + \left( \frac{n^2}{2} + n + \frac{5}{12} + \frac{z^2}{2} + (n+1)z \right) \ln n \\ & - \frac{3n^2}{4} - n - nz - \ln A + \frac{1}{12} + O(n^{-1}), \end{aligned}$$

where  $\Gamma$  is the familiar (Euler's) gamma function and  $A$  is called the Glaisher-Kinkelin constant defined by

$$(4) \quad \ln A = \lim_{n \rightarrow \infty} \left\{ \ln \left( \prod_{k=1}^n k^k \right) - \left( \frac{n^2}{2} + \frac{n}{2} + \frac{1}{12} \right) \ln n + \frac{n^2}{4} \right\},$$

the numerical value of  $A$  being  $1.28242713 \dots$ .

The Glaisher-Kinkelin constant  $A$  can be expressed as follows (see [36]):

$$A = \lim_{n \rightarrow \infty} n^{-n^2/2 - n/2 - 1/12} e^{n^2/4} \prod_{k=1}^n k^k,$$

$$\frac{e^{1/12}}{A} = \lim_{n \rightarrow \infty} \frac{G(n+1)}{n^{n^2/2 - 1/12} (2\pi)^{n/2} e^{-3n^2/4}},$$

and (see [25, p. 129, Eq. (3.22)])

$$(5) \quad A = e^{\frac{1}{12} - \zeta'(-1)} = (2\pi)^{1/12} [e^{\gamma\pi^2/6 - \zeta'(2)}]^{1/(2\pi^2)},$$

where  $\zeta'(z)$  is the derivative of the Riemann zeta function  $\zeta(z)$  (see [30]).

The Glaisher-Kinkelin constant  $A$  has drawn attention in many works (see, for example, [16, 20, 25, 29, 30]; see also [4]). One section in the book by Finch [37, pp. 135–138] was devoted to introduce this Glaisher-Kinkelin constant  $A$ . Moreover, the Glaisher-Kinkelin constant  $A$  plays an important role in the study of the Barnes  $G$ -function (see, for details, [66, Section 1.4]).

The following integral representation for the remainder  $R_N(z)$  of the explicit expression for the Barnes  $G$ -function was established by Ferreira and López [36, Theorem 1].

**Theorem 1.** *An integral representation for the remainder  $R_N(z)$  in the following explicit expression for the Barnes  $G$ -function:*

$$\begin{aligned} \ln G(z+1) &= \frac{1}{4}z^2 + z \ln \Gamma(z+1) - \left( \frac{1}{2}z^2 + \frac{1}{2}z + \frac{1}{12} \right) \ln z - \ln A \\ &+ \sum_{k=1}^{N-1} \frac{B_{2k+2}}{2k(2k+1)(2k+2)z^{2k}} + R_N(z) \quad (N \in \mathbb{N}; |\operatorname{Arg}(z)| < \pi) \end{aligned}$$

where  $B_{2k+2}$  are the Bernoulli numbers, is given for  $\Re(z) > 0$  by

$$R_N(z) = \int_0^\infty \left( \frac{t}{e^t - 1} - \sum_{k=0}^{2N} \frac{B_k}{k!} t^k \right) \frac{e^{-zt}}{t^3} dt \quad (\Re(z) > 0).$$

Estimates for  $|R_N(z)|$  were also found by Ferreira and López [36], showing that the expansion is indeed an asymptotic expansion of  $\ln G(z+1)$  in the sectors of the complex plane cut along the negative real axis. Pedersen [59, Theorem 1.1] proved that, for any  $N \geq 1$ , the function  $x \mapsto (-1)^N R_N(x)$  is completely monotonic on  $(0, \infty)$ . Other asymptotic relations (avoiding the  $\ln \Gamma$  term) was given by Ruijsenaars [62] and investigated by Pedersen [60], Koumandos [41], and by Koumandos and Pedersen [42]. Some upper and lower bounds for the double gamma function were derived in terms of the gamma, psi and polygamma functions

(see [8, 9, 14, 22]). Chen [15] and Mortici [54] established several inequalities and asymptotic expansions for  $\ln A$  in (4). Chen and Lin [20] and Chen [16] presented a class of asymptotic expansions related to the Glaisher-Kinkelin constant  $A$  and the Barnes  $G$ -function. Recently, Chen and Srivastava [23] presented a number of potentially useful properties of the Barnes  $G$ -function. The properties considered here include, for example, its integral representation, complete monotonicity, and continued-fraction approximation. We also derive continued-fraction approximations of the Glaisher-Kinkelin constant  $A$  and the Choi-Srivastava constants  $B$  and  $C$ , which are analogous to the Glaisher-Kinkelin constant  $A$  and are given by (see [29, p.102] and [30])

$$\ln B = \lim_{n \rightarrow \infty} \left\{ \sum_{k=1}^n k^2 \ln k - \left( \frac{n^3}{3} + \frac{n^2}{2} + \frac{n}{6} \right) \ln n + \frac{n^3}{9} - \frac{n}{12} \right\}$$

and

$$\ln C = \lim_{n \rightarrow \infty} \left\{ \sum_{k=1}^n k^3 \ln k - \left( \frac{n^4}{4} + \frac{n^3}{2} + \frac{n^2}{4} - \frac{1}{120} \right) \ln n + \frac{n^4}{16} - \frac{n^2}{12} \right\}$$

for which the approximate numerical values are given by

$$B = 1.03091675 \dots \quad \text{and} \quad C = 0.97955746 \dots$$

As  $x \rightarrow \infty$ , the Stirling formula for the Barnes  $G$ -function can be found as follows (see [65, p. 26]):

$$(6) \quad \ln G(x+1) = \frac{x}{2} \ln(2\pi) - \ln A + \frac{1}{12} - \frac{3x^2}{4} + \left( \frac{x^2}{2} - \frac{1}{12} \right) \ln x + O(x^{-1}).$$

Chen [17] applied the formula (6) to produce the following complete asymptotic expansion:

$$(7) \quad \begin{aligned} \ln G(x+1) \sim & \frac{x}{2} \ln(2\pi) - \ln A + \frac{1}{12} - \frac{3x^2}{4} + \left( \frac{x^2}{2} - \frac{1}{12} \right) \ln x \\ & - \frac{1}{240x^2} + \frac{1}{1008x^4} - \frac{1}{1440x^6} + \frac{1}{1056x^8} - \frac{691}{327600x^{10}} + \dots \end{aligned}$$

and derived a recurrence relation for determining the coefficients of  $1/x^j$  ( $j \geq 2$ ) occurring in the expansion (7).

Just as the expression of the Glaisher-Kinkelin constant  $A$  in (5), the Choi-Srivastava constants  $B$  and  $C$  are also known to be expressible in terms of special values of the derivative of the Riemann zeta function  $\zeta(s)$  as follows (see [30] and [31, Eq. (1.9)]):

$$(8) \quad \ln B = -\zeta'(-2) \quad \text{and} \quad \ln C = -\frac{11}{720} - \zeta'(-3).$$

Chen [15] established the asymptotic expansions related to the Glaisher-Kinkelin constant  $A$  and the Choi-Srivastava constants  $B$  and  $C$ . Mortici [54] also dealt with the same problem. Cheng and Chen [24] and Chen and Choi [18] established some novel asymptotic expansions of the Glaisher-Kinkelin constant  $A$  and the Choi-Srivastava constants  $B$  and  $C$ . Also, by using the Bernoulli numbers  $B_n$ , Chen [15] established the asymptotic expansions related to the constants  $A$ ,  $B$  and  $C$ .

Recently, Chen [17] derived a recurrence relation for determining the coefficients of each asymptotic expansion related to the constants  $A$ ,  $B$  and  $C$ , without using the Bernoulli numbers  $B_n$ . More precisely, Chen [17] proved the following results:

$$\begin{aligned} \sum_{k=1}^n k \ln k - \left( \frac{n^2}{2} + \frac{n}{2} + \frac{1}{12} \right) \ln n + \frac{n^2}{4} \\ \sim \ln A + \frac{1}{720n^2} - \frac{1}{5040n^4} + \frac{1}{10080n^6} - \frac{1}{9504n^8} + \dots, \end{aligned}$$

$$\begin{aligned} \sum_{k=1}^n k^2 \ln k - \left( \frac{n^3}{3} + \frac{n^2}{2} + \frac{n}{6} \right) \ln n + \frac{n^3}{9} - \frac{n}{12} \\ \sim \ln B - \frac{1}{360n} + \frac{1}{7560n^3} - \frac{1}{25200n^5} + \frac{1}{33264n^7} - \frac{691}{16216200n^9} + \dots \end{aligned}$$

and

$$\begin{aligned} \sum_{k=1}^n k^3 \ln k - \left( \frac{n^4}{4} + \frac{n^3}{2} + \frac{n^2}{4} - \frac{1}{120} \right) \ln n + \frac{n^4}{16} - \frac{n^2}{12} \\ \sim \ln C - \frac{1}{5040n^2} + \frac{1}{33600n^4} - \frac{1}{66528n^6} + \frac{691}{43243200n^8} - \frac{1}{34320n^{10}} + \dots, \end{aligned}$$

as  $n \rightarrow \infty$  in each case.

In our present investigation, we introduce two new mathematical constants  $D$  and  $E$ , which are analogous to the constants  $A$ ,  $B$  and  $C$ , defined by

$$(9) \quad \ln D = \lim_{n \rightarrow \infty} \left\{ \sum_{k=1}^n k^4 \ln k - \left( \frac{1}{5} n^5 + \frac{1}{2} n^4 + \frac{1}{3} n^3 - \frac{1}{30} n \right) \ln n + \frac{1}{25} n^5 - \frac{1}{12} n^3 + \frac{13}{360} n \right\}$$

and

$$(10) \quad \ln E = \lim_{n \rightarrow \infty} \left\{ \sum_{k=1}^n k^5 \ln k - \left( \frac{1}{6} n^6 + \frac{1}{2} n^5 + \frac{5}{12} n^4 - \frac{1}{12} n^2 + \frac{1}{252} \right) \ln n + \frac{1}{36} n^6 - \frac{1}{12} n^4 + \frac{47}{720} n^2 \right\},$$

respectively, approximate numerical values of  $D$  and  $E$  being given by

$$D = 0.99204797 \dots \quad \text{and} \quad E = 1.00968038 \dots .$$

We also derive a recurrence relation for determining the coefficients of each of the asymptotic expansions, which are related to the constants  $D$  and  $E$ , without using the Bernoulli numbers  $B_n$ .

**Remark.** We thank a referee for drawing our attention toward the related developments in [1, 11, 27]. The mathematical constants  $D$  and  $E$ , which we have introduced and studied in this work, were studied in a generalized form which includes the Glaisher-Kinkelin constant  $A$  as well as the Choi-srivastava constants  $B$  and  $C$  (see, for details, [1, 11]; see also [25, p. 131, Eq. (4.10)] for the corrected form of [1, p. 198, Eq. (20)] as well as [27]).

## 2. ASYMPTOTIC EXPANSIONS FOR THE WALLIS SEQUENCE

Theorem 2 below provides a recurrence relation for determining the coefficients of  $n^{-j}$  in the expansion (1), without the help of the Bernoulli numbers  $B_n$ .

**Theorem 2.** *As  $n \rightarrow \infty$ , the following asymptotic expansion holds true:*

$$W_n \sim \frac{\pi}{2} \exp\left(\sum_{k=1}^{\infty} \frac{\nu_k}{n^k}\right),$$

where the coefficients  $\nu_j$  are given by the recurrence relation given by

$$(11) \quad \nu_1 = -\frac{1}{4} \quad \text{and} \quad \nu_k = \frac{(-1)^k}{k} \left\{ \frac{3^{k+1} - 2^{k+2} + 1}{2^{k+1}(k+1)} - \sum_{j=1}^{k-1} \nu_j (-1)^j \binom{k}{k-j+1} \right\}$$

for  $k \geq 2$ .

*Proof.* Upon setting

$$P_n = \ln\left(\frac{2}{\pi} W_n\right) \quad \text{and} \quad Q_n = \sum_{k=1}^{\infty} \frac{\nu_k}{n^k},$$

we can let  $P_n \sim Q_n$  and

$$\Delta P_n := P_{n+1} - P_n \sim Q_{n+1} - Q_n =: \Delta Q_n$$

as  $n \rightarrow \infty$ , where  $\nu_k$  are real numbers to be determined.

By making use of the fact that

$$W_n = \frac{\pi}{2} \cdot \frac{1}{n + \frac{1}{2}} \left( \frac{\Gamma(n+1)}{\Gamma(n + \frac{1}{2})} \right)^2 = \frac{\pi}{2} \cdot \frac{\Gamma(n+1)^2}{\Gamma(n + \frac{1}{2})\Gamma(n + \frac{3}{2})},$$

we obtain

$$\begin{aligned} \Delta P_n &= 2 \ln \left( 1 + \frac{1}{n} \right) - \ln \left( 1 + \frac{1}{2n} \right) - \ln \left( 1 + \frac{3}{2n} \right) \\ (12) \quad &= \sum_{k=2}^{\infty} \frac{(-1)^k (3^k - 2^{k+1} + 1)}{2^k k} n^{-k}. \end{aligned}$$

Also, by direct computation, we get

$$\begin{aligned} \sum_{k=1}^{\infty} \frac{\nu_k}{(n+1)^k} &= \sum_{k=1}^{\infty} \frac{\nu_k}{n^k} \left( 1 + \frac{1}{n} \right)^{-k} = \sum_{k=1}^{\infty} \frac{\nu_k}{n^k} \sum_{j=0}^{\infty} \binom{-k}{j} \frac{1}{n^j} \\ &= \sum_{k=1}^{\infty} \frac{\nu_k}{n^k} \sum_{j=0}^{\infty} (-1)^j \binom{k+j-1}{j} \frac{1}{n^j} \\ (13) \quad &= \sum_{k=1}^{\infty} \left\{ \sum_{j=1}^k \nu_j (-1)^{k-j} \binom{k-1}{k-j} \right\} n^{-k}. \end{aligned}$$

We thus find that

$$(14) \quad \Delta Q_n = \sum_{k=2}^{\infty} \left\{ \sum_{j=1}^k \nu_j (-1)^{k-j} \binom{k-1}{k-j} - \nu_k \right\} n^{-k}.$$

Now, upon equating the coefficients of  $n^{-k}$  on the right-hand sides of (12) and (14), we have

$$\begin{aligned} \frac{(-1)^k (3^k - 2^{k+1} + 1)}{2^k k} &= \sum_{j=1}^k \nu_j (-1)^{k-j} \binom{k-1}{k-j} - \nu_k \\ &= \sum_{j=1}^{k-1} \nu_j (-1)^{k-j} \binom{k-1}{k-j} \end{aligned}$$

and

$$\begin{aligned} \frac{3^k - 2^{k+1} + 1}{2^k k} &= \sum_{j=1}^{k-1} \nu_j (-1)^j \binom{k-1}{k-j} \\ &= \sum_{j=1}^{k-2} \nu_j (-1)^j \binom{k-1}{k-j} + (-1)^{k-1} (k-1) \nu_{k-1} \quad (k \geq 2), \end{aligned}$$



where (and elsewhere in this paper) an empty sum is understood to be zero.

For  $k = 2$ , we obtain  $\nu_1 = -\frac{1}{4}$ . Also, for  $k \geq 3$ , we have

$$\nu_{k-1} = \frac{(-1)^{k-1}}{k-1} \left\{ \frac{3^k - 2^{k+1} + 1}{2^k k} - \sum_{j=1}^{k-2} \nu_j (-1)^j \binom{k-1}{k-j} \right\},$$

which can be written precisely as (11). The proof of Theorem 2 is thus completed.  $\square$

Our next result (Theorem 3) provides a recurrence relation for determining the coefficients of  $n^{-j}$  in the expansion (3) without the coefficients  $\nu_j$ .

**Theorem 3.** *As  $n \rightarrow \infty$ , the following asymptotic expansion holds true:*

$$W_n \sim \frac{\pi}{2} \sum_{k=0}^{\infty} \frac{\mu_k}{n^k},$$

where the coefficients  $\mu_k$  are given by the recurrence relation:

$$(15) \quad \mu_0 = 1, \quad \mu_1 = -\frac{1}{4}, \quad \text{and} \quad \mu_k = -\frac{1}{k} \left\{ \frac{2k^2 - 2k + 1}{4} \mu_{k-1} + \sum_{j=0}^{k-2} \mu_j \binom{k}{k-j+1} \right\}$$

for  $k \geq 2$ .

*Proof.* We first set

$$U_n = \frac{2}{\pi} W_n \quad \text{and} \quad V_n = \sum_{k=0}^{\infty} \frac{\mu_k}{n^k},$$

where  $\mu_0 = 1$ . We can then let  $U_n \sim V_n$  and

$$\frac{U_n}{U_{n-1}} \sim \frac{V_n}{V_{n-1}} \quad (n \rightarrow \infty),$$

where  $\mu_k$  are real numbers to be determined as follows:

$$\frac{4n^2}{4n^2 - 1} \sim \frac{\sum_{k=0}^{\infty} \frac{\mu_k}{n^k}}{\sum_{k=0}^{\infty} \frac{\mu_k}{(n-1)^k}},$$

which yields

$$(16) \quad \sum_{k=0}^{\infty} \frac{\mu_k}{(n-1)^k} \sim \left(1 - \frac{1}{4n^2}\right) \sum_{k=0}^{\infty} \frac{\mu_k}{n^k} = \mu_0 + \frac{\mu_1}{n} + \sum_{k=2}^{\infty} \left(\mu_k - \frac{\mu_{k-2}}{4}\right) \frac{1}{n^k}.$$

Now, by direct computation, we get

$$\begin{aligned}
 \sum_{k=0}^{\infty} \frac{\mu_k}{(n-1)^k} &= \sum_{k=0}^{\infty} \frac{\mu_k}{n^k} \left(1 - \frac{1}{n}\right)^{-k} = \sum_{k=0}^{\infty} \frac{\mu_k}{n^k} \sum_{j=0}^{\infty} (-1)^j \binom{-k}{j} \frac{1}{n^j} \\
 &= \sum_{k=0}^{\infty} \frac{\mu_k}{n^k} \sum_{j=0}^{\infty} \binom{k+j-1}{j} \frac{1}{n^j} = \sum_{k=0}^{\infty} \sum_{j=0}^k \mu_j \binom{k-1}{k-j} \frac{1}{n^k} \\
 (17) \quad &= \mu_0 + \frac{\mu_1}{n} + \sum_{k=2}^{\infty} \sum_{j=0}^k \mu_j \binom{k-1}{k-j} \frac{1}{n^k}.
 \end{aligned}$$

Thus, upon equating the coefficients of  $n^{-k}$  on the right-hand sides of (16) and (17), we find that

$$\mu_k - \frac{\mu_{k-2}}{4} = \sum_{j=0}^k \mu_j \binom{k-1}{k-j} \quad (k \geq 2)$$

and

$$\begin{aligned}
 -\frac{\mu_{k-2}}{4} &= \sum_{j=0}^{k-1} \mu_j \binom{k-1}{k-j} = \sum_{j=0}^{k-2} \mu_j \binom{k-1}{k-j} + (k-1)\mu_{k-1} \\
 &= \sum_{j=0}^{k-3} \mu_j \binom{k-1}{k-j} + \frac{(k-1)(k-2)}{2} \mu_{k-2} + (k-1)\mu_{k-1}.
 \end{aligned}$$

For  $k = 2$ , we obtain  $\mu_1 = -\frac{1}{4}$ . Also, for  $k \geq 3$ , we have

$$\mu_{k-1} = -\frac{1}{k-1} \left\{ \left( \frac{1}{4} + \frac{(k-1)(k-2)}{2} \right) \mu_{k-2} + \sum_{j=0}^{k-3} \mu_j \binom{k-1}{k-j} \right\},$$

which can be written precisely as (15). This completes our proof of Theorem 3.  $\square$

### 3. ASYMPTOTIC EXPANSIONS RELATED TO THE CONSTANTS $D$ AND $E$

In this section, we first recall the Euler-Maclaurin summation formula as follows (see, for example, [39, p. 318]; see also [34]):

$$(18) \quad \sum_{k=1}^n f(k) \sim C_0 + \int_a^n f(x) dx + \frac{1}{2} f(n) + \sum_{r=1}^{\infty} \frac{B_{2r}}{(2r)!} f^{(2r-1)}(n),$$

where  $C_0$  is an arbitrary constant to be determined in each special case and the  $B_{2k}$  are the Bernoulli numbers. Indeed, if we set

$$f(x) = x^4 \ln x \quad \text{and} \quad f(x) = x^5 \ln x$$

in (18) with  $a = 1$ , we are led to (9) and (10), respectively.

As in the cases of  $\ln A$  in (5),  $\ln B$  and  $\ln C$  in (8), we can also express  $\ln D$  and  $\ln E$  as special cases of  $\zeta'(s)$ . In this connection, by using the Euler-Maclaurin summation formula (18), we can obtain a number of analytical representations of  $\zeta(s)$ , such as the known result recorded by Hardy [39, p. 333]:

$$\zeta(s) = \lim_{n \rightarrow \infty} \left\{ \sum_{k=1}^n k^{-s} - \frac{n^{1-s}}{1-s} - \frac{1}{2} n^{-s} \right\} \quad (\Re(z) > -1),$$

$$\zeta(s) = \lim_{n \rightarrow \infty} \left\{ \sum_{k=1}^n k^{-s} - \frac{n^{1-s}}{1-s} - \frac{1}{2} n^{-s} + \frac{1}{12} sn^{-s-1} \right\} \quad (\Re(z) > -3),$$

$$\begin{aligned} \zeta(s) = \lim_{n \rightarrow \infty} \left\{ \sum_{k=1}^n k^{-s} - \frac{n^{1-s}}{1-s} - \frac{1}{2} n^{-s} + \frac{1}{12} sn^{-s-1} \right. \\ \left. - \frac{1}{720} s(s+1)(s+2)n^{-s-3} \right\} \quad (\Re(z) > -5) \end{aligned} \tag{19}$$

and

$$\begin{aligned} \zeta(s) = \lim_{n \rightarrow \infty} \left\{ \sum_{k=1}^n k^{-s} - \frac{n^{1-s}}{1-s} - \frac{1}{2} n^{-s} + \frac{1}{12} sn^{-s-1} - \frac{1}{720} s(s+1)(s+2)n^{-s-3} \right. \\ \left. + \frac{1}{30240} s(s+1)(s+2)(s+3)(s+4)n^{-s-5} \right\} \quad (\Re(z) > -7). \end{aligned} \tag{20}$$

By first differentiating both sides of (19) with respect to  $s$  and then setting  $s = -4$ , we obtain

$$\begin{aligned} -\zeta'(-4) = \lim_{n \rightarrow \infty} \left\{ \sum_{k=1}^n k^4 \ln k - \left( \frac{1}{5} n^5 + \frac{1}{2} n^4 + \frac{1}{3} n^3 - \frac{1}{30} n \right) \ln n \right. \\ \left. + \frac{1}{25} n^5 - \frac{1}{12} n^3 + \frac{13}{360} n \right\}, \end{aligned}$$

which, when compared with (9), yields

$$\ln D = -\zeta'(-4).$$

Also, by first differentiating both sides of (20) with respect to  $s$  and then

setting  $s = -5$ , we obtain

$$\frac{137}{15120} - \zeta'(-5) = \lim_{n \rightarrow \infty} \left\{ \sum_{k=1}^n k^5 \ln k - \left( \frac{1}{6} n^6 + \frac{1}{2} n^5 + \frac{5}{12} n^4 - \frac{1}{12} n^2 + \frac{1}{252} \right) \ln n + \frac{1}{36} n^6 - \frac{1}{12} n^4 + \frac{47}{720} n^2 \right\},$$

which, when compared with (10), yields

$$\ln E = \frac{137}{15120} - \zeta'(-5).$$

Without using the Bernoulli numbers, Theorems 4 and 5 provide a recurrence relation for determining the coefficients of each asymptotic expansion related to the above-defined constants  $D$  and  $E$  in (9) and (10), respectively.

**Theorem 4.** *As  $n \rightarrow \infty$ , the following asymptotic expansion holds true:*

$$\begin{aligned} \sum_{k=1}^n k^4 \ln k - \left( \frac{1}{5} n^5 + \frac{1}{2} n^4 + \frac{1}{3} n^3 - \frac{1}{30} n \right) \ln n + \frac{1}{25} n^5 - \frac{1}{12} n^3 + \frac{13}{360} n \\ \sim \ln D + \sum_{k=1}^{\infty} \frac{d_k}{n^k}, \end{aligned}$$

where the coefficients  $d_k$  are given by the recurrence relation:

$$(21) \quad d_k = \frac{(-1)^{k+1}(k+14)}{30(k+2)(k+4)(k+5)(k+6)} - \frac{1}{k} \sum_{j=1}^{k-1} d_j (-1)^{k-j} \binom{k}{k-j+1} \quad (k \geq 2)$$

together with  $d_1 = \frac{1}{1260}$ . In terms of the constant  $D$  defined by (9), it is asserted that

$$\begin{aligned} \sum_{k=1}^n k^4 \ln k - \left( \frac{1}{5} n^5 + \frac{1}{2} n^4 + \frac{1}{3} n^3 - \frac{1}{30} n \right) \ln n + \frac{1}{25} n^5 - \frac{1}{12} n^3 + \frac{13}{360} n \\ \sim \ln D + \frac{1}{1260n} - \frac{1}{25200n^3} + \frac{1}{83160n^5} - \frac{691}{75675600n^7} + \cdots \end{aligned}$$

*Proof.* If we set

$$\begin{aligned} X_n = \sum_{k=1}^n k^4 \ln k - \left( \frac{1}{5} n^5 + \frac{1}{2} n^4 + \frac{1}{3} n^3 - \frac{1}{30} n \right) \ln n \\ + \frac{1}{25} n^5 - \frac{1}{12} n^3 + \frac{13}{360} n - \ln D \end{aligned}$$

and

$$Y_n = \sum_{k=1}^{\infty} \frac{d_k}{n^k},$$

then we can let  $X_n \sim Y_n$  and

$$\Delta X_n := X_{n+1} - X_n \sim Y_{n+1} - Y_n =: \Delta Y_n \quad (n \rightarrow \infty),$$

where  $d_k$  are real numbers to be determined. Indeed, after some elementary transformations, we find that

$$\begin{aligned} \Delta X_n &= \left( -\frac{1}{5} n^5 - \frac{1}{2} n^4 - \frac{1}{3} n^3 + \frac{1}{30} n \right) \ln \left( 1 + \frac{1}{n} \right) \\ &\quad + \frac{2}{5} n^3 - \frac{1}{20} n + \frac{1}{5} n^4 + \frac{3}{20} n^2 - \frac{13}{1800} \\ (22) \quad &= \sum_{k=2}^{\infty} \frac{(-1)^{k-1} (k-1)(k+13)}{30(k+1)(k+3)(k+4)(k+5)} \frac{1}{n^k}. \end{aligned}$$

Thus, by using (13), we obtain

$$(23) \quad \Delta V_n = \sum_{k=2}^{\infty} \left\{ \sum_{j=1}^k d_j (-1)^{k-j} \binom{k-1}{k-j} - d_k \right\} \frac{1}{n^k}.$$

Upon equating the coefficients of  $n^{-k}$  on the right-hand sides of (22) and (23), we get

$$\begin{aligned} \frac{(-1)^{k-1} (k-1)(k+13)}{30(k+1)(k+3)(k+4)(k+5)} &= \sum_{j=1}^k d_j (-1)^{k-j} \binom{k-1}{k-j} - d_k \\ &= \sum_{j=1}^{k-1} d_j (-1)^{k-j} \binom{k-1}{k-j} \\ &= \sum_{j=1}^{k-2} d_j (-1)^{k-j} \binom{k-1}{k-j} - (k-1)d_{k-1} \quad (k \geq 2), \end{aligned}$$

which, for  $k = 2$ , yields  $d_1 = \frac{1}{1260}$ . Moreover, for  $k \geq 3$ , we have

$$d_{k-1} = \frac{(-1)^k (k+13)}{30(k+1)(k+3)(k+4)(k+5)} + \frac{1}{k-1} \sum_{j=1}^{k-2} d_j (-1)^{k-j} \binom{k-1}{k-j},$$

which can be written as (21). The proof of Theorem 4 is now complete.  $\square$

**Theorem 5.** *As  $n \rightarrow \infty$ , the following asymptotic expansion holds true:*

$$\sum_{k=1}^n k^5 \ln k - \left( \frac{1}{6} n^6 + \frac{1}{2} n^5 + \frac{5}{12} n^4 - \frac{1}{12} n^2 + \frac{1}{252} \right) \ln n + \frac{1}{36} n^6 - \frac{1}{12} n^4 + \frac{47}{720} n^2 \\ \sim \ln E + \sum_{k=1}^{\infty} \frac{e_k}{n^k},$$

where the coefficients  $e_k$  are given by the recurrence relation:

$$(24) \quad e_k = \frac{(-1)^k (k-1)(141 + 22k + k^2)}{252(k+1)(k+3)(k+5)(k+6)(k+7)} - \frac{1}{k} \sum_{j=1}^{k-1} e_j (-1)^{k-j} \binom{k}{k-j+1} \quad (k \geq 2)$$

together with  $e_1 = 0$ . Furthermore, in terms of the constant  $E$  defined by (10), it is asserted that

$$\sum_{k=1}^n k^5 \ln k - \left( \frac{1}{6} n^6 + \frac{1}{2} n^5 + \frac{5}{12} n^4 - \frac{1}{12} n^2 + \frac{1}{252} \right) \ln n + \frac{1}{36} n^6 - \frac{1}{12} n^4 + \frac{47}{720} n^2 \\ \sim \ln E + \frac{1}{10080n^2} - \frac{1}{66528n^4} + \frac{691}{90810720n^6} - \frac{1}{123552n^8} + \dots$$

*Proof.* Upon setting

$$I_n = \sum_{k=1}^n k^5 \ln k - \left( \frac{1}{6} n^6 + \frac{1}{2} n^5 + \frac{5}{12} n^4 - \frac{1}{12} n^2 + \frac{1}{252} \right) \ln n \\ + \frac{1}{36} n^6 - \frac{1}{12} n^4 + \frac{47}{720} n^2 - \ln E$$

and

$$J_n = \sum_{k=1}^{\infty} \frac{e_k}{n^k},$$

we can let  $I_n \sim J_n$  and

$$\Delta I_n := I_{n+1} - I_n \sim J_{n+1} - J_n =: \Delta J_n \quad (n \rightarrow \infty),$$

where  $e_k$  are real numbers to be determined. In fact, after some elementary transformations, we obtain

$$(25) \quad \Delta I_n = \left( -\frac{1}{6} n^6 - \frac{1}{2} n^5 - \frac{5}{12} n^4 + \frac{1}{12} n^2 - \frac{1}{252} \right) \ln \left( 1 + \frac{1}{n} \right) \\ + \frac{1}{6} n^5 + \frac{5}{12} n^4 + \frac{2}{9} n^3 - \frac{1}{12} n^2 - \frac{13}{360} n + \frac{7}{720} \\ = \sum_{k=3}^{\infty} \frac{(-1)^k (k-1)(k-2)(k^2 + 20k + 120)}{252k(k+2)(k+4)(k+5)(k+6)} \frac{1}{n^k}$$

and

$$(26) \quad \Delta J_n = -\frac{e_1}{n^2} + \sum_{k=3}^{\infty} \left\{ \sum_{j=1}^k e_j (-1)^{k-j} \binom{k-1}{k-j} - e_k \right\} \frac{1}{n^k}.$$

Now, by equating the coefficients of  $n^{-k}$  on the right-hand sides of (25) and (26), we find that  $e_1 = 0$  and that

$$\begin{aligned} \frac{(-1)^k (k-1)(k-2)(k^2+20k+120)}{252k(k+2)(k+4)(k+5)(k+6)} &= \sum_{j=1}^k e_j (-1)^{k-j} \binom{k-1}{k-j} - e_k \\ &= \sum_{j=1}^{k-1} e_j (-1)^{k-j} \binom{k-1}{k-j} \\ &= \sum_{j=1}^{k-2} e_j (-1)^{k-j} \binom{k-1}{k-j} - (k-1)e_{k-1} \end{aligned}$$

and

$$e_{k-1} = \frac{1}{k-1} \sum_{j=1}^{k-2} e_j (-1)^{k-j} \binom{k-1}{k-j} - \frac{(-1)^k (k-2)(k^2+20k+120)}{252k(k+2)(k+4)(k+5)(k+6)} \quad (k \geq 3),$$

which can be written precisely as (24). This evidently completes our proof of Theorem 5.  $\square$

#### 4. AN OPEN PROBLEM

As the Euler-Mascheroni constant  $\gamma$  is involved with the classical gamma function  $\Gamma$ , the constants  $A$ ,  $B$  and  $C$  have appeared naturally in the theory of the multiple gamma functions  $\Gamma_n$  (see, for example, [66, Section 1.4]) and play their respective roles as described in ([65, p. 39, p. 247], [28, p. 523, Eq. (2.50)] and [25]).

$$(27) \quad \int_0^{\frac{1}{2}} \ln \Gamma(t+1) dt = -\frac{1}{2} - \frac{7}{24} \ln 2 + \frac{1}{4} \ln \pi + \frac{3}{2} \ln A,$$

$$(28) \quad \int_0^{\frac{1}{2}} \ln G(t+1) dt = \frac{1}{24} (\ln 2 + 1) + \frac{1}{16} \ln \pi - \frac{1}{4} \ln A - \frac{7}{4} \ln B$$

and

$$(29) \quad \int_0^{\frac{3}{2}} \ln \Gamma_3(t+2) dt = -\frac{259}{768} - \frac{29}{1920} \ln 2 + \frac{9}{16} \ln \pi - \frac{15}{16} \ln A - \frac{5}{4} \ln B + \frac{15}{16} \ln C,$$

where  $\Gamma_3$  is the triple gamma function (see [66, p. 58]).

In view of (27), (28) and (29), we propose the following open problem.

**Open Problem.** Let  $\alpha$  and  $\beta$  be two given positive numbers. Determine the constants  $a_j \equiv a_j(\alpha, \beta)$  and  $b_j \equiv b_j(\alpha, \beta)$  such that

$$\int_0^\alpha \ln \Gamma_4(t + \beta) dt = a_1 + a_2 \ln 2 + a_3 \ln \pi + a_4 \ln A + a_5 \ln B \\ + a_6 \ln C + a_7 \ln D$$

and

$$\int_0^\alpha \ln \Gamma_5(t + \beta) dt = b_1 + b_2 \ln 2 + b_3 \ln \pi + b_4 \ln A + b_5 \ln B \\ + b_6 \ln C + b_7 \ln D + b_8 \ln E.$$

## 5. CONCLUDING REMARKS AND OBSERVATIONS

In our present investigation, without using the Bernoulli numbers  $B_n$ , we have established several asymptotic expansions and a recurrence relation for determining the coefficients of each asymptotic expansion associated with the Wallis sequence  $W_n$ , defined by

$$W_n := \prod_{k=1}^n \frac{4k^2}{4k^2 - 1},$$

and the constants  $D$  and  $E$ , which are analogous to the Glaisher-Kinkelin constant  $A$  and the Choi-Srivastava constants  $B$  and  $C$ . We have also pointed out the relevant connections of the formulas and results, which we have considered in this article, with various known or new results.

**Acknowledgements.** The authors thank the referee for helpful comments. This work was supported by the Key Science Research Project in the Universities of Henan Province (20B110007) and the Fundamental Research Funds for the Universities of the Henan Province (NSFRF210446).

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(Received 14. 04. 2022.)

(Revised 25. 08. 2022.)

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